A Quasi-Analytical Method for Non-iterative Computation of Nonlinear Controls

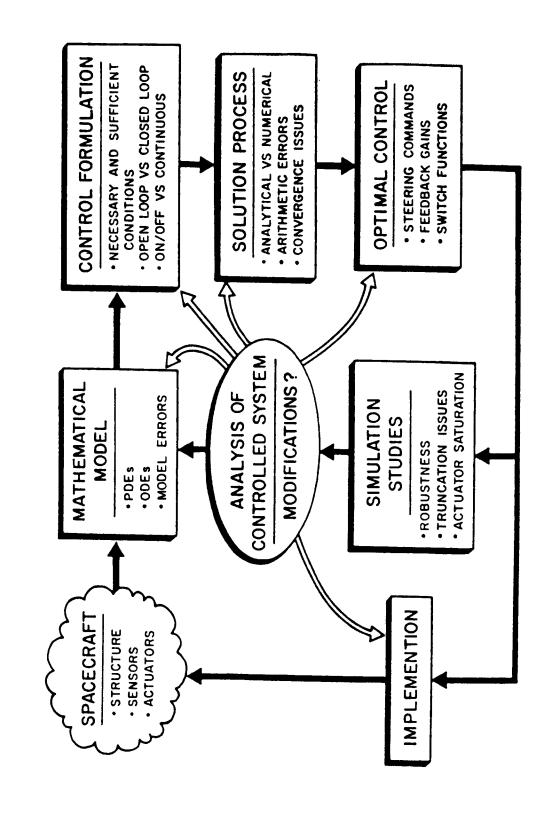
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and

J. D. Turner Cambridge Research Associates workshop on Structural Dynamics and Control
INTERACTION FOR FLEXIBLE STRUCTURES
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Coupling of Spacecraft Structural Modeling with Dynamics/Controls Analysis, Design, and Implementation*



Junkins, J. L. and Turner, J. D., Optimal Spacecraft Rotational Maneuvers, Elsevier, 1986. the above figure is from our book:

PRELIMINARIES

Consider a dynamical system described by

$$\dot{z} = Fz + Du + \epsilon g(z, u, t)$$

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where

z is an nxl state vector

u is an mxl control vector

F & D are constant matrices

z, u, and the nonlinear terms g(z,u,t) are continuous & differentiable

We seek an optimal control $u^*(t)$ and corresponding optimal trajectory $z^*(t)$ which minimize the quadratic performance measure

$$J = 112 \left[z^{T} S z \right]_{t_{f}} + 112 \int_{S}^{t_{f}} (z^{T} Q z + u^{T} R u) dt \tag{2}$$

The necessary conditions involve the Hamiltonian H(z, u, p, t)

$$H = 1/2 (z^{T}Qz + u^{T}Ru) + p^{T}(Fz + Du + \epsilon g);$$
 (3)

hese are:
$$\frac{\partial H}{\partial u} = 0$$
, $\frac{\partial H}{\partial p} = \dot{z}$, $-\frac{\partial H}{\partial z} = \dot{p}$

plus boundary conditions: $z(0) = z_0$, $p(t_f) = Sz(t_f)$ or $z(t_f) = z_f$.

Two Point Boundary Value Problem

Pontryagin Necessary Conditions:

$$\dot{z} = Fz + Du + \epsilon g$$
, $z(0) = z_0$
 $\dot{p} = -Qz - F^T p - \epsilon \left[\frac{\partial g}{\partial z} \right]^T$, $p(t_f) = Sz(t_f)$, for $z(t_f)$ "free"

 $0 = Ru + D^{T}p$, ... from which the optimal control is $u^{*} = -R^{-1}D^{T}p$.

The state/co-state coupled system can be written as a 2n order system:

$$\dot{x} = A x + \epsilon h(x,t)$$
where
$$x^{T} = [z^{T} \ p^{T}], A = \begin{bmatrix} F & -DR^{T}D^{T} \\ \\ \\ -Q & -F^{T} \end{bmatrix}, h(x,t) = \begin{bmatrix} g \\ \\ -\frac{\partial g^{T}}{\partial z} \end{bmatrix}$$
(4)

through use of a perturbation method & "quasi-analytical" integration. >>> are often expensive due to initial ignorance of a "good starting estimate" (of Due to the nonlinear terms in h(x,t), exact analytical solutions of Eq. (4) are $p(t_o)$ or $p(t_f)$) required for reliable convergence. We seek to avoid iteration most often impossible. A variety of iterative techniques are available; they

The Asymptotic Expansion of the Necessary Conditions

We seek a power series solution of the usual form

$$x(t) = x_0(t) + \epsilon x_1(t) + \epsilon^2 x_2(t) + \cdots + \epsilon^k x_k(t) + \cdots$$
 (5)

Substitution of the power series into the state/co-state system of Eq. (4), and equating like powers of ←, leads to the <u>sequence</u> of linear systems:

$$\dot{x}_{o} = A x_{o}$$

$$\dot{x}_{I} = A x_{I} + g_{I}(t, x_{o}(t))$$

$$- x_{I}(t)$$
(6)

$$x_k = A x_k + g_k(t, x_0(t), x_1(t), ..., x_{k-1}(t)) - x_k(t)$$

with the boundary conditions

$$\left. egin{aligned} x_o(o) = \left\langle \begin{array}{c} z(t_f) \\ y, x_o(t_g) = \left\langle \begin{array}{c} z(t_f) \\ y, x_o(t_g) \end{array} \right\rangle \end{array} \right. , \quad \left. \begin{array}{c} z(t_f) \\ y, \dots, x_k(t_f) = \left\langle \begin{array}{c} 0 \\ p_k(t_f) \end{array} \right\rangle \right. \end{array}$$

where, at least formally, the sequence of solutions is given by

$$x_k(t) = e^{At} \left[x_k(0) + \int_0^t e^{-A\tau} g_k(\tau, x_0(\tau), x_I(\tau), \dots, x_{k-I}(\tau)) d\tau \right], \, k = 1, 2, 3, \dots$$
 (7)

But... how do we make efficient algorithms? Does convergence occur in the "real world"? Can the above be implemented in a way which automates the algebra usually associated with perturbation methods? What about secular terms? Does this approach apply to systems of non-trivial dimensions & "messy" nonlinear terms? We have made some progress in answering these questions.

Van Loan's Identity for Integrating Forced Linear Systems

Conside

$$\dot{x}_k = A x_k + g_k, \quad x_{k(t)} = e^{At} x_{k(0)} + e^{At} \int_{-1}^{1} e^{-A\tau} g_{k(\tau)} d\tau \quad , k = 1, 2, \dots \quad (8)$$

For the special case that U(t) can be represented as Fourier series, the Fourier series can be rewritten as a matrix exponential

$$g_k(t) = b_{ok} + \sum_{r=1}^{N} b_{rk} \cos(\omega_r t) + a_{rk} \sin(\omega_r t) = G_k e^{\Omega t} c$$
 (9)

where

 $oldsymbol{b}_{ok}$, $oldsymbol{b}_{rk}$, $oldsymbol{a}_{rk}$, are 2n x I vectors of Fourier coefficients

$$G_k = [b_o b_1 a_1 b_2 a_2 \dots b_r a_r \dots b_N a_N]_k$$
, a 2n x (2N + 1) constant matrix

$$c = \{1 \ 1 \ 0 \ 1 \ 0 \ 1 \ 0 \dots 1 \ 0\}^T$$
, $a_{2N \times 1}$ selection vector

$$\Omega_r = \begin{bmatrix} 0 & -\omega_r \\ \omega_r & 0 \end{bmatrix}, \quad \omega_r = r(2\pi/(t_f - t_o))$$

$$\Omega = \operatorname{diag} [0, \Omega_I, \Omega_2, \dots, \Omega_r, \dots, \Omega_N]$$

Substituting Eq. (9) into Eq. (8),

$$x_k(t) = e^{At} x_k(0) + [\int_{0}^{t} e^{-At} G_k e^{-At}]c = e^{At} x_k(0) + [v_k]c$$

Van Loan has established the interesting & useful identity which permits computation of the

forced response using a matrix exponential (via, for example, Ward's Pade' algorithm):
$$\begin{bmatrix} A & G_t \\ O & \Omega \end{bmatrix}_t = \begin{bmatrix} e & At & | & \psi_k \\ & - & | & - & | \\ & 0 & | & e & \Omega t \end{bmatrix}$$
(10)

For large N, we can use superposition & keep the order of the matrix exponentials small thus the response to a relatively arbitrary $g_k(t)$ can be calculated via matrix exponentials.

Control Rate Smoothing & State Vector Augmentation

We choose to minimize

$$J = 112 \{ z(t_f)^T S z(t_f) + u(t_f)^T S_0 u(t_f) + \dot{u}(t_f)^T S_1 \dot{u}(t_f) \}$$

$$+ 112 \int_0^t \{ z^T Q z + u^T R_0 u + \dot{u}^T R_1 \dot{u} + \ddot{u}^T R_2 \ddot{u} \} dt$$

Subject to: $\dot{z} = Az + Du$. This can be converted to standard form via the definitions:

$$\tilde{z} = \begin{pmatrix} z \\ u \\ \dot{u} \end{pmatrix}, \quad \tilde{A} = \begin{bmatrix} A & D & 0 \\ 0 & 0 & I \\ 0 & 0 & 0 \end{bmatrix}, \quad \tilde{D} = \begin{bmatrix} 0 \\ 0 \\ I \end{bmatrix}, \quad \tilde{Q} = block diag[Q, R_0, R_I]$$

So we can equivalently minimize

$$J = 1/2 \ \tilde{z}(t_f)^T \tilde{S} \tilde{z}(t_f) + 1/2 \ \int_{\Omega}^{t_f} \{ \tilde{z}^T \tilde{Q} \tilde{z} + \tilde{u}^T \tilde{R} \tilde{u} \} dt$$

profiles to decrease excitation of the poorly modeled higher frequency modes. form as those developed in the foregoing. Penalizing the control derivatives Subject to: $\tilde{z} = \tilde{A}\tilde{z} + \tilde{D}\tilde{u}$. The necessary conditions have the identical has been found most constructive in frequency-shaping the torque

Case 1 Optimal Detumble/Attitude Aquisition

STATE DYNAMICS

Euler (quaternion) parameters

$$\begin{vmatrix} \dot{\beta}_0 \\ \dot{\dot{\beta}}_1 \\ \dot{\dot{\beta}}_2 \\ \dot{\dot{\beta}}_3 \\ \dot{\dot{\beta}}_3 \\ \dot{\dot{\beta}}_3 \\ \end{pmatrix} = \frac{1}{2} \begin{bmatrix} 0 & -\omega_1 & -\omega_2 & -\omega_3 \\ \omega_1 & 0 & \omega_3 & -\omega_2 \\ \omega_2 & -\omega_3 & 0 & \omega_1 \\ \omega_2 & -\omega_3 & 0 & \omega_1 \\ \omega_3 & \omega_2 & -\omega_1 & 0 \end{bmatrix} \begin{pmatrix} \beta_0 \\ \beta_2 \\ \beta_3 \\ \end{pmatrix}$$

$$\int_{0}^{\infty} \operatorname{or} \dot{\mathbf{g}} = \frac{1}{2} (\omega) \, \mathbf{g}$$

Euler's Equations

$$\begin{pmatrix} \dot{\omega}_1 \\ \dot{\omega}_2 \\ \dot{\omega}_2 \\ \dot{\omega}_3 \end{pmatrix} = \begin{pmatrix} -I_1 & \omega_2 \omega_3 + & u_1/I_1 \\ -I_2 & \omega_3 \omega_1 + & u_2/I_2 \\ -I_3 & \omega_1 \omega_2 + & u_3/I_3 \end{pmatrix} , \quad I_2 = \frac{(I_1 - I_2)/I_1}{I_2 - I_1/I_3}$$

or $\dot{\omega} = f(\omega, u)$

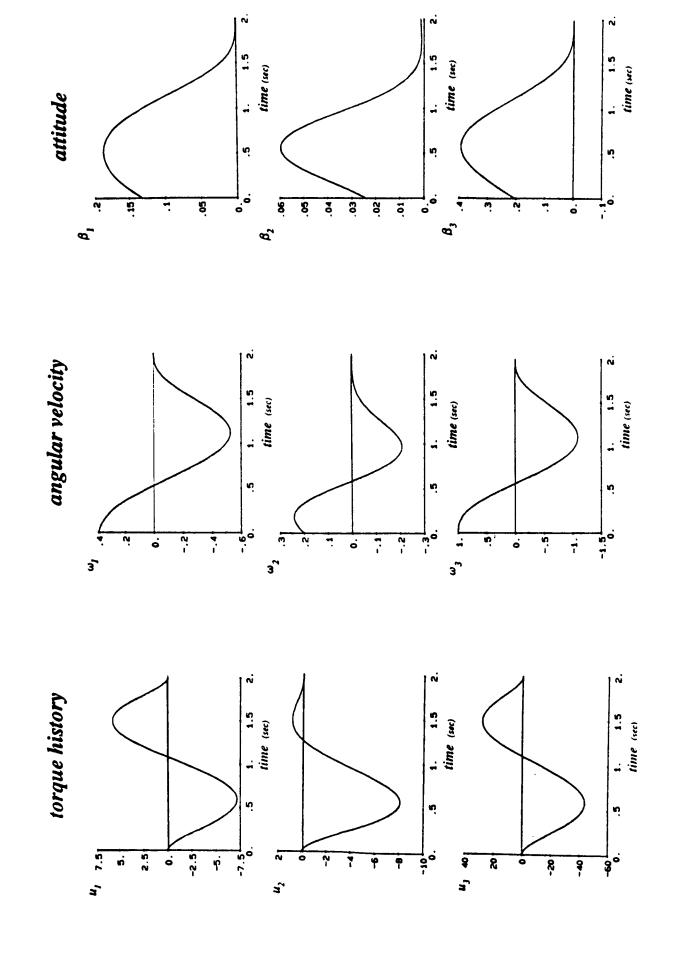
(0) =
$$\begin{cases} .9699665 \\ .1318887 \\ .0238626 \\ .2029798 \end{cases}$$
, $\omega(0) = \begin{cases} .4 \text{ r/s} \\ .2 \\ 1.0 \end{cases}$ $B(t_f) = \begin{cases} 1 \\ 0 \\ 0 \end{cases}$, $\omega(t_f) = \begin{cases} 0 \\ 0 \\ 0 \end{cases}$

Case 1 Numerical Results for the TPBVP Solution

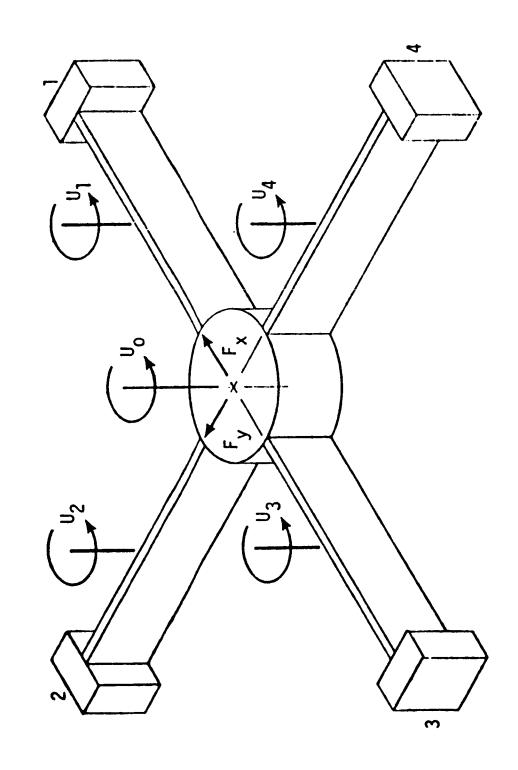
FINAL STATE ERRORS

	LINEAR SO	LUTION	FIRST ORDER	ROER	SECOND ORDER	ORDER
ΔβΩ	,01999		16000.		7×10^{-7}	
081	08232	× 10°	99800*-	°2.	.00058	° 50° ≈
Δ82	.18033		00944		.00094	
- ۵۵۶	.01734		00412		00034	
ζωγ	.01914		01525		.00109	
- Δω ₂	.43150		00292		.00295	
	.00461		00071		.00015	

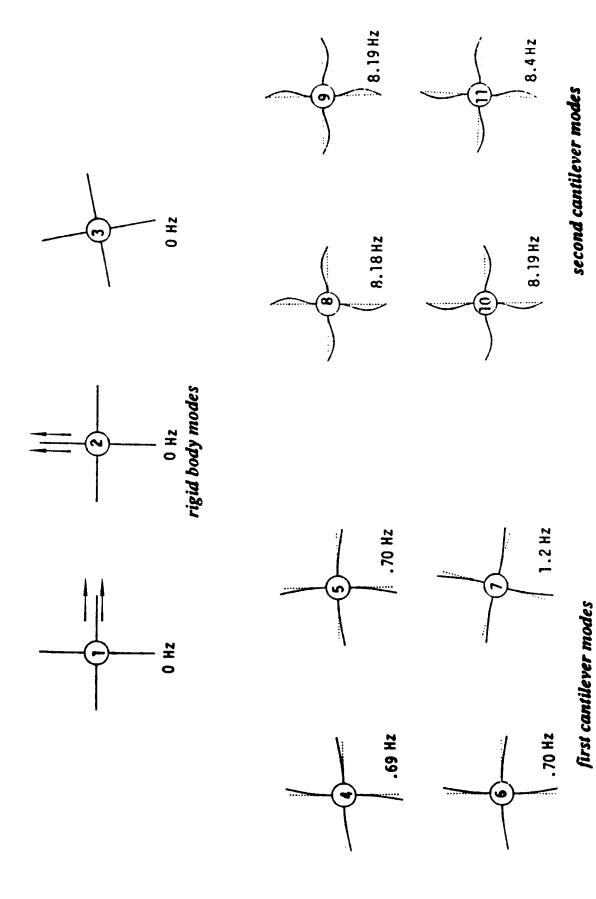
Case I Optimal Detumble/Attitude Acquisition Maneuver



The Draper/RPL Slewing Experimental Configuration



Draper/RPL Configuration: First Eleven in-plane Vibration Modes



Case 2 Numerical Results for the TPBVP Solution (small angle flexible body maneuver)

FINAL STATE ERRORS

	LINEAR SOLUTION	FIRST ORDER	SECOND ORDER
+	-0.828E-5	0.276E-4	0.184E-5
0	0.104E-2	0.667E-5	-0.878E-6
->	0.668E-3	-0.270E-4	0.130E-5
3	-0.450E-3	0.150E-3	0.342E-5
8	0.317E-2	0.293E-4	-0.474E-5
3	0.597E-3	-0.396E-5	0.135E-6
η ₁ (in plane)	-0.960E-3	-0.120E-5	-0.217E-8
σ, (out-of-plane)	0.317E-2	0.293E-4	-0.474E-5

anti-symmetric 'anti-symmetric out-of-plane in-plane 1. 1.5 time (100) 1. 1.5 time (sec) deflection ເນ 'n F .0021 -.02 .02 -.01 .001 -.001 9. ö Case 2 Optimal Maneuver with Vibration Suppression/Arrest i... time (sec) 1. 1.5 time (sec) 1. 1.5 time (sec) attitude 'n ĸ B00. 900. .004 .002 0 6 .05 .03 .02 9 -02 9 90. 9 'n 1. 1.5 time (sec) 1. 1.5 time (sec) angular velocity 1. 1.5 time (sec) 'n 3 3 - . - . -.02 -.04 -.08 -.002 -.004 -.06 -.006 -.01 -.02 -.03 -.04 .05 0 1. 1.5 time (100) 5. 1. 1.5 time (sec) time (sec) torque history 7 93 .05 'n 0

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Concluding Remarks

A novel optimal control solution process has been developed for a general class of nonlinear dynamical systems

The method combines control theory, perturbation methods, and Van Loan's recent matrix exponential results

matrix exponentials (using Ward's Pade algorithm) and All controlled response integrations are accomplished via recursions developed herein A variety of applications support the practical utility of this method; flexible body dynamical systems of order >40 have been solved nonlinear rigid body optimal maneuvers are routinely solved;

The method fails occasionally due to poor convergence of the perturbation expansion or numerical difficulties associated with computing the matrix exponential

for semi-automation; no initial guess is required, and it usually The method is attractive because it appears to be a good candidate converges at 2nd or 3rd order in minutes of machine time

REFERENCES

- Slewing Guidance Feedback Spacecraft," AIAA 82-4131, Journal of Control, Vol. 5, No. 3, May-June, 1982, p. 318. "Comment on Optimal J.L., Junkins, Flexible ٦.
- Solution for the Linear Tracking Problems," Paper No. 83-374, presented at the AAS/AIAA Astrodynamics Conference, Lake Turner, J.D., Chun, H.M. and Juang, Jer-Nan, "Optimal Slewing a Closed Form Maneuvers for Flexible Spacecraft Using Placid, New York, August 22-25, 1983. 2
- Nayfeh, Ali Hasan, and Mook, Dean T., Nonlinear Oscillations, John Wiley and Sons, New York, 1979. . m
- Exponential, "IEEE Transactions on Auto Control, Vol. AC-23, No. 3, June, 1978, pp. 395-404. the Van Loan, C.F., "Computing Integrals Involving 4.
- Ward, R.C., "Numerical Computation of the Matrix Exponential with Accuracy Estimate," SIAM Journal of Numerical Analysis, Vol. 14, No. 4, September, 1977, pp. 600-610. ۍ ک
- Kirk, Donald E., Optimal Control Theory An Introduction, Prentice Hall, Inc., Englewood Cliffs, NJ, 1970. 9
- R.C., A Perturbation Approach to Control Flexible Thesis, Engineering Mechanics, of Maneuvers Rotational/Translational Blacksburg, VA, 1985. M.S. Thompson, 7
- Junkins, J.L. and Thompson, R.C., "An Asymptotic Perturbation Method for Nonlinear Optimal Control Problems," Paper No. AAS 85-364, presented at the AAS/AIAA Astrodynamics Specialist Conference, Vail, Colorado, August, 12-15, 1985. . &
- Junkins, J.L. and Turner, J.D., Optimal Spacecraft Rotational Maneuvers, Elsevier, Amsterdam, 1986. 6